

Optical Dephasing and Acoustic Plasmon Undamping in Highly Excited Semiconductors

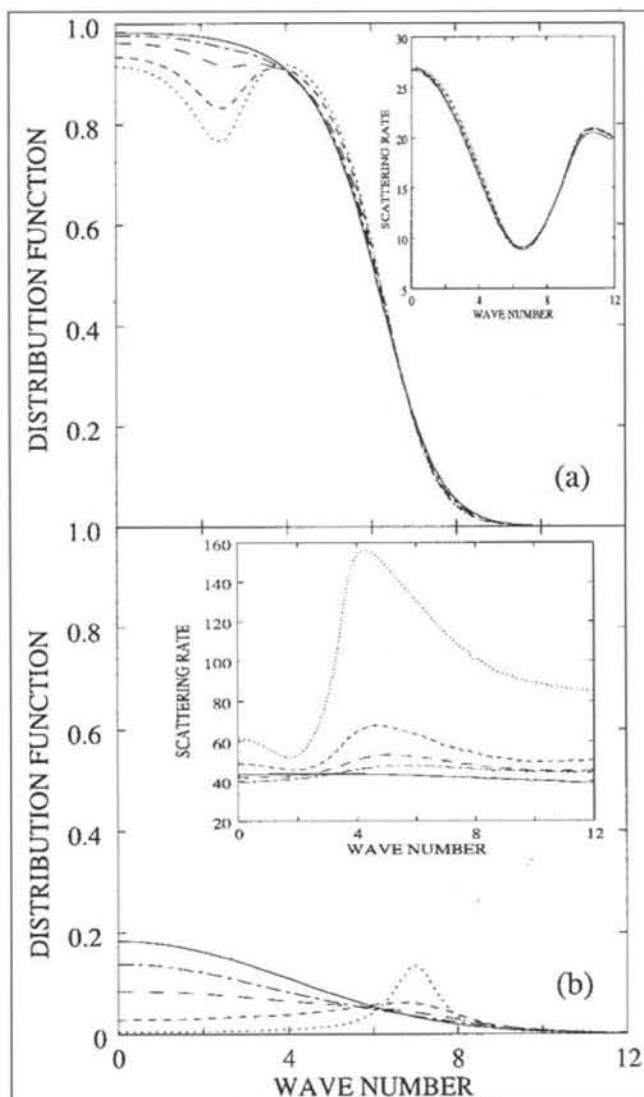
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Optical dephasing, *i.e.*, the decay of the polarization field in a semiconductor, is a direct consequence of electron and hole scattering. Under high excitation conditions or in a semiconductor amplifier/laser, carrier-carrier scattering is often the dominant relaxation mechanism, which also leads to energy-level broadening and dynamical screening of the Coulomb interaction potential. The theoretical analysis of carrier-carrier scattering is based on the quantum Boltzmann equation. Even though this equation is well known in the many-body literature,¹ its solution for high excitation conditions and nonequilibrium carrier distributions is a substantial challenge.

To study nonequilibrium carrier relaxation and optical dephasing, we performed direct numerical integrations of the time-dependent Boltzmann equation where the dynamical screening of the Coulomb interaction potential was consistently treated at the same level as the scattering probabilities.^{2,3} The solution of the quantum Boltzmann equation yields the dynamic changes of the electron and hole distribution functions. These changes can be expressed in terms of scattering rates, which determine the average time an electron or hole remains in a quantum state. In generalization of the well-known relation between optical dephasing and excited state lifetime in atomic systems, the sum of the "lifetimes" of the semiconductor electrons and holes in a state determine the dephasing of that state.^{2,3}

We studied the carrier dynamics and optical dephasing for an optical semiconductor amplifier, *i.e.*, an inverted semiconductor where the carrier density is sufficiently high to allow a spectral region of optical gain. The amplification of an incident light beam leads to a frequency dependent depletion of the inversion, *i.e.*, to a "hole" in the distribution function. Studying the relaxation of such a "kinetic hole," we obtained carrier relaxation and optical dephasing times of the order of 50 fsec for carrier densities, typical for amplifier or laser operation conditions in GaAs.² The figure shows an example of the dynamic evolution of the electron distribution function. The corresponding scattering rates are essentially time independent. The momentum-dependence of these scattering rates reflects the phase-space restriction near the Fermi edge, which is strongly smeared out in this case.

A very different situation occurs for femtosecond excitation of passive semiconductors, energetically high in the region of optical interband absorption. For the generated nonequilibrium plasma, we find significantly reduced relaxation times, which can be as short as 10 fsec.³ In Part B of the figure, we plot the relaxation of the nonequilibrium electron-hole plasma distribution. The dynamic scattering rates (inset) exhibit a very strong momentum dependence with pronounced peaks corresponding to dephasing times around



Numerical solutions of the electron-hole Boltzmann equation using the dynamically screened Coulomb potential in random phase approximation. The figure shows only the electron distributions at different times; the full electron-hole relaxation dynamics can be found in Refs. 2 and 3. (a) Relaxation of initially disturbed Fermi distribution functions for density $n = 3 \times 10^{18} \text{ cm}^{-3}$ and temperature $T = 300 \text{ K}$. Shown is the electron distribution as function of the carrier momentum in units of the exciton Bohr radius (a_B). The initial ($t = 0$) distribution function is plotted as a dotted line. Consecutive times are $t = 21$ fsec (short-dashed), $t = 75$ fsec (long-dashed), and $t = 147$ fsec (dash-dotted), and the final time $t = 796$ fsec (solid). (b) Relaxation of initial nonequilibrium distribution for density $n = 4.6 \times 10^{17} \text{ cm}^{-3}$. The initial ($t = 0$) curve is plotted as a dotted line. Consecutive times are $t = 70$ fsec (short-dashed line), $t = 185$ fsec (long-dashed line), $t = 355$ fsec (dash-dotted line), and the final time $t = 1590$ fsec (solid line). The insets show the corresponding scattering rates in ps^{-1} .

10 fsec. We find that such extremely high scattering rates occur only for carrier configurations that allow the undamping of the acoustic plasmon resonance,³ which is strongly Landau damped under equilibrium conditions. Our analysis shows that the acoustic plasmon undamping for nonequilibrium distributions causes the enhancement of the scattering matrix element leading to ultrafast relaxation and optical dephasing.

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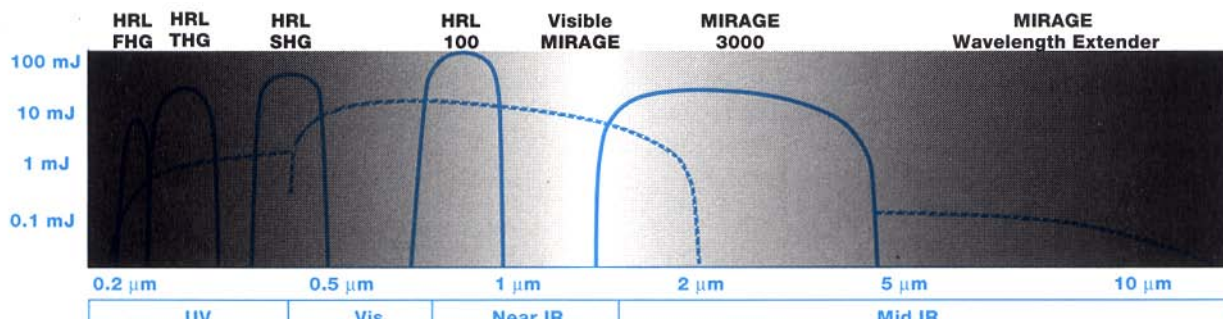


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